

The effective Einstein's equation and the GW E-M tensor ①

$$R_{\mu\nu} = \frac{8\pi G}{c^4} \left(T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \right) \quad (1)$$

Once we consider $g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$ we have

$$R_{\mu\nu} = \bar{R}_{\mu\nu} + R_{\mu\nu}^{(1)} + R_{\mu\nu}^{(2)} + \dots \quad (2)$$

Now we calculate the effective tensor :

$$\langle R_{\mu\nu} \rangle = \bar{R}_{\mu\nu} + \langle R_{\mu\nu}^{(2)L} \rangle \quad (3) \text{ hence :}$$

$$\bar{R}_{\mu\nu} = - \langle R_{\mu\nu}^{(2)L} \rangle + \frac{8\pi G}{c^4} \langle T_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} T \rangle \quad (4)$$

$$\text{We define } t_{\mu\nu} = - \frac{c^4}{8\pi G} \langle R_{\mu\nu}^{(2)} - \frac{1}{2} \bar{g}_{\mu\nu} R^{(2)} \rangle \quad (5) \text{ where}$$

$$R^{(2)} = \bar{g}^{\mu\nu} R_{\mu\nu}^{(2)} \rightarrow t = \bar{g}^{\mu\nu} t_{\mu\nu} = - \frac{c^4}{8\pi G} \langle R^{(2)} - \frac{1}{2} 4 R^{(2)} \rangle = \\ = \frac{c^4}{8\pi G} \langle R^{(2)} \rangle \quad (6)$$

Replacing (6) into (5) we have

$$t_{\mu\nu} = - \frac{c^4}{8\pi G} \langle R_{\mu\nu}^{(2)} \rangle + \bar{g}_{\mu\nu} t \quad \text{i.e.:}$$

$$\langle R_{\mu\nu}^{(2)} \rangle = - \frac{8\pi G}{c^4} (t_{\mu\nu} - \bar{g}_{\mu\nu} t) \quad (7)$$

Replacing (7) in (4):

(2)

$$\bar{R}_{\mu\nu} = \frac{8\pi G}{c^4} \left[(t_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} t) + (\bar{T}_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T}) \right] \quad (8)$$

Once again, multiplying by $\bar{g}_{\mu\nu}$ we have:

$$\bar{R} = \frac{8\pi G}{c^4} [-t - \bar{T}] \quad \text{hence}$$

$$\bar{R}_{\mu\nu} = \frac{8\pi G}{c^4} (\bar{T}_{\mu\nu} + t_{\mu\nu}) + \frac{1}{2} \bar{g}_{\mu\nu} \frac{8\pi G}{c^4} (t + \bar{T}) \quad \text{i.e.}$$

$$\boxed{\bar{R}_{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{R} = \frac{8\pi G}{c^4} (\bar{T}_{\mu\nu} + t_{\mu\nu})} \quad (9)$$

This is the low-frequency part of the Einstein's equation, otherwise called the coarse-grained equation.

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The E-M tensor for GW is then

$$t_{\mu\nu} = -\frac{c^4}{8\pi G} \left[\langle R_{\mu\nu}^{(2)} \rangle - \frac{1}{2} \bar{g}_{\mu\nu} \langle R^{(2)} \rangle \right] \quad (5)$$

After long calculations and considering the TT gauge we have

$$\text{top. } \langle R_{\mu\nu}^{(2)} \rangle = -\frac{1}{4} \langle \partial_\mu h_{\alpha\beta} \partial_\nu h^{\alpha\beta} \rangle \quad (10)$$

First consequence:

$$\begin{aligned}
 R^{(2)} &= \langle g^{\mu\nu} R_{\mu\nu}^{(2)} \rangle = -\frac{1}{4} \langle \partial^\nu h_{\alpha\beta} \partial_\nu h^{\alpha\beta} \rangle = \\
 &= -\frac{1}{4} \left\{ \int h_{\alpha\beta} \partial_\nu h^{\alpha\beta} \cdot n^\nu dS \right. \\
 &\quad \left. \begin{array}{l} \text{surface} \downarrow \\ = 0 \end{array} \right. \neq \underbrace{\langle \partial^\nu \partial_\nu h_{\alpha\beta} h^{\alpha\beta} \rangle}_{\square h_{\alpha\beta} = 0} = \\
 &= 0
 \end{aligned} \tag{11}$$

$$\text{So, } t_{\mu\nu} = \frac{c^4}{32\pi G} \langle \partial_\mu h_{\alpha\beta} \partial_\nu h^{\alpha\beta} \rangle \tag{12}$$

$$\begin{aligned}
 t_{00} &= \frac{c^2}{32\pi G} \langle \dot{h}_{\alpha\beta}^{\text{TT}} \cdot \dot{h}^{\alpha\beta} \rangle = \\
 &= \frac{c^2}{32\pi G} \langle \dot{h}_+^2 + \dot{h}_\times^2 \rangle \tag{13}
 \end{aligned}$$